

COULFG: COULOMB AND BESSEL FUNCTIONS AND THEIR DERIVATIVES, FOR REAL ARGUMENTS, BY STEED'S METHOD

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Received 15 December 1981

PROGRAM SUMMARY

Title of program: COULFG: Coulomb, Bessel Functions

Catalogue number: ABNK

Program obtainable from: CPC Program Library, Queen's University of Belfast, N. Ireland (see application form in this issue)

Computer: IBM 370/165 and AS/7000; *Installation:* Daresbury Laboratory, Warrington, Lancs.

Operating system: OS/360 GI compiler and HX compiler (level 2.2.1)

Programming language used: ASA FORTRAN

High speed storage required: 180 Kbytes

No. of bits in a word: 32

Overlay structure: none

Peripherals used: card reader, printer

No. of cards in combined program and test deck: 432

Card punching code: EBCDIC

Keywords: Klein-Gordon, Coulomb for real angular momentum, recurrence relations, Schrödinger, Bessel, spherical Bessel, continued fraction, reactions, scattering, heavy ion, nuclear, molecular, atomic, pionic, kaonic, exotic atoms, scattering states

Nature of physical problem

Coulomb interaction of charged particles in spherical coordinates (Coulomb function) and uncharged particles (spherical Bessel functions), and general problems in cylindrical coordinates yielding Bessel-function solutions can be solved with the program. COULFG computes the Coulomb wavefunctions $F_\lambda(\eta, x)$, $G_\lambda(\eta, x)$, $F'_\lambda(\eta, x)$ and $G'_\lambda(\eta, x)$ for a range of λ -val-

ues in integer steps, $l \geq \lambda \geq m > -1$, for real $x > 0$ and real η , $10^4 \geq |\eta|$. Values of the maximum angular momentum, l , of several thousand can be treated. When the functions have oscillating character the absolute accuracies are about 10^{-14} but this is only limited by the word length used; extended precision variables can yield absolute accuracies of 10^{-30} . COULFG will generate spherical Bessel functions and cylindrical Bessel functions for a wide range of integer-spaced real orders.

Method of solution

An enhanced version of Steed's method, used previously for integer λ in subroutine RCWFN [1] is adopted. The more recent subroutine KLEIN [2] for a single λ value also is similar, the additional feature being the stable recurrence relations for a range of λ values.

Restriction on the complexity of the problem

It is well known that the method used loses accuracy as $x < x_m$ (the turning point for the minimum order required) and eventually, when $G_m \gtrsim 10^6$, a JWKB approximate solution is adopted which is accurate to $\lesssim 1\%$ as a rule. Several output variables signal that this has occurred.

Features of the program

A considerably revised version is presented of the CPC program 'RCWFN' (catalogue ABPC) used to compute the Coulomb functions $F_L(\eta, x)$, $G_L(\eta, x)$ and their x -derivatives over a range of integer L -values ≥ 0 . The new program, COULFG, calculates $F_\lambda(\eta, x)$, $G_\lambda(\eta, x)$, $F'_\lambda(\eta, x)$ and $G'_\lambda(\eta, x)$ when λ is real ($\lambda > -1$) and for a similar range in the (η, x) plane as before, $10^4 \geq x > 0$, $10^4 \geq |\eta|$. Integer-spaced λ -values are obtained by a suitable combination of stable recurrence relations. Subroutine COULFG will return, furthermore, for a range of integer-spaced orders;

- a) Spherical Bessel functions $j_\lambda(x)$, $y_\lambda(x)$, $j'_\lambda(x)$, $y'_\lambda(x)$, and
- b) Cylindrical Bessel functions $J_\mu(x)$, $Y_\mu(x)$, $J'_\mu(x)$, $Y'_\mu(x)$, where both λ , η can be real or integral. Values of order in excess of 1000 have been tested.

A "mode" option is provided so that, for each case, arrays of the regular and irregular functions and their derivatives, or just the two functions, or the regular function alone ($F_\lambda, j_\lambda, J_\mu$) are calculated and stored. Both core and execution time are saved by this technique.

In the region of x where the functions oscillate, i.e. $x \geq x_\lambda$ (the turning point for the λ th partial wave) the relative accuracy of the functions for IBM machines (REAL*8) as programmed, is $\approx 10^{-14}$ – 10^{-16} and for CDC machines (single precision) is $\approx 10^{-12}$ – 10^{-14} . With no change in the code except for the accuracy parameter the accuracy can be increased to $\approx 10^{-30}$ by using the AUTODOUBLE facility on an extended-precision IBM compiler. The accuracy decreases in a predictable way as

x decreases below x_λ ; a measure of this decreased accuracy is provided by the program. JWKB approximations are provided when x is sufficiently smaller than x_λ .

Typical running time

The test deck ran in $2\frac{1}{2}$ s on the GI compiler and the HX compiler, one half of which demonstrates error conditions.

References

- [1] A.R. Barnett, D.H. Feng, J.W. Steed and L.J.B. Goldfarb, *Comput. Phys. Commun.* 8 (1974) 377.
- [2] A.R. Barnett, *Comput. Phys. Commun.* 24 (1981) 141.

LONG WRITE-UP

1. Introduction

This is the third paper in the current series which discusses programs deriving from Steed's method of calculating Coulomb wavefunctions. The previous papers, to be referred to as I, II, dealt with an overview of the algorithms and their family relationships [1], and with KLEIN [2], a one-shot program for a fixed real angular momentum quantum number, λ , real η and real x (i.e. positive energy). KLEIN returns the values of the regular solution, $F_\lambda(\eta, x)$, the irregular solution, $G_\lambda(\eta, x)$ and their x -derivatives to the same accuracy. The present paper describes COULFG, which is designed to extend the methods of KLEIN to a range of integer-spaced λ -values, $m \leq \lambda \leq l$, where $l - m \equiv 0 \pmod{1}$ and $m > -1$, and this is equivalent for integer $\lambda = L$ to the original subroutine RCWFN of Barnett et al. [3]. Hence it provides Coulomb wavefunction solutions to the non-relativistic Schrödinger equation for real $\eta \geq 0$. For $m = l$ the solutions to the relativistic Schrödinger equation, i.e. the Klein–Gordon equation, are obtained, by setting (see I, II for notation)

$$\lambda = \left[\left(L + \frac{1}{2} \right)^2 - Z^2 \alpha^2 \right]^{1/2} - \frac{1}{2}, \quad (1)$$

and the results of KLEIN are reproduced. By the use of similar expressions for λ , given in II, scattering-state solutions with positive energy of both the (relativistic) Dirac equation, and the Klein–Gordon equation in N -dimensional space,

can be obtained. In addition, both Bessel functions and spherical Bessel functions (regular and irregular) over a wide range of orders can be obtained readily from COULFG. The accuracies of all functions, when x is in the oscillating region, are about 10^{-14} – 10^{-16} in the standard (REAL*8) IBM program using 56-bit words; on a single precision CDC compiler with 48-bit words the accuracy becomes 10^{-12} – 10^{-14} . By changing one card (line 107) and using an extended-precision IBM compiler with the AUTODOUBLE facility the accuracy can be squared ($\approx 10^{-30}$). Subsequent papers will deal with programs specifically for Bessel functions, for the real Gamma function and for the Airy functions.

The present program COULFG supercedes RCWFN [3] in its capabilities but is identical in principle and for identical parameters will produce identical results, with the exception that for x -values below the turning point

$$x_\lambda = \eta + \left[\eta^2 + \lambda(\lambda + 1) \right]^{1/2} \quad (2)$$

for the minimum $\lambda = m$, the earlier program provided an integration method for G and G' , which has been removed. If x is sufficiently smaller than x_m , that is for $F \lesssim 10^{-6}$ and $G \gtrsim 10^6$, a JWKB approximation, see II, is employed which in general provides results to better than 1%. The specific methods of the standard programs of Bardin et al. [4] treat the small- x case directly.

The subroutine COULFG includes KLEIN (in

effect) and it includes the improvements made in RCWFF [5]; thus it incorporates a 'mode' parameter such that when $\text{MODE} = 1$ then F_λ , G_λ , F'_λ and G'_λ are used in the calculation and the arrays filled, when $\text{MODE} = 2$ only the arrays of the functions F_λ , G_λ are used and filled, while when $\text{MODE} = 3$ only the array F_λ is used and has results stored in it. This technique economises on core storage and time when some of the quantities are not required. Thus, in part of a typical application [6], values of F_λ from $m = 0$ to $l = 800$ were required but not the irregular solution or the derivatives. Thus the array declaration statement `DIMENSION FC(801), GC(1), FCP(1), GCP(1)` was adequate in the calling program.

An extensive description of the algorithm, also called COULFG, is given in I and full details of the Coulomb functions for real λ appeared in II and in related papers which examine Steed's method [3] in detail for accuracy [7], and range and limitations [8], and comprehensive references to the extensive literature on Coulomb functions are given there. The most recent summary [9] is that of Kölbig in 1972, while Fullerton [10] has compiled a large bibliography of calculations of special functions which includes the Coulomb functions but not Klein-Gordon ones.

The relativistic Coulomb equation, solved by COULFG for real $x > 0$, real $\lambda > -1$ and real η , including $\eta = 0$, is

$$d^2U/dx^2 + [1 - 2\eta/x - \lambda(\lambda + 1)/x^2] U_\lambda = 0. \quad (3)$$

The next section details the relationship to the Bessel functions and the method of solution, section 3 contains program notes and the last section describes the test output and program listing.

2. Subroutine COULFG

2.1. The Coulomb functions

The programming closely follows the algorithm specified in I (with some errata corrected [11]) and the details in II. Familiarity with these two references will be assumed and the range of λ -values will be taken as $m \leq \lambda \leq l$, with $l - m$ an integer. First the ratio $f_l = F'_l(\eta, x)/F_l(\eta, x)$ is obtained for

the maximum λ -value (XLMAX) required, by means of the continued fraction CF1. The relative accuracy is set by the internal parameter ACCUR, which is 10^{-16} in the listing. Then f_l is set to $s\beta$ and F'_l to $s\beta f_l$, where $\beta = 10^{-30}$ and $s = \text{sign } F_l$. The recurrence relations [1],

$$\begin{aligned} F_{\lambda-1} &= (S_\lambda F_\lambda + F'_\lambda)/R_\lambda, \\ F'_{\lambda-1} &= S_\lambda F'_{\lambda-1} - R_\lambda F_\lambda, \end{aligned} \quad (4)$$

are used to find $F_m(\eta, x)$ and $F'_m(\eta, x)$, with F'_λ not being stored if $\text{MODE} = 2, 3$. Their ratio, F'_m/F_m , is f_m and it is combined with $p_m + iq_m$ obtained by means of the continued fraction CF2 (see I, II) to obtain F_m , G_m , F'_m and G'_m at the minimum λ -value required (XLMIN). The regular function is, for example,

$$F_m(\eta, x) = \text{sign}(F_m) [(f_m - p_m)^2/q_m + q_m]^{-1/2}, \quad (5)$$

(the expression following (25) in I is corrected here [11]), and the others are

$$F'_m = f_m F_m; \quad G_m = \gamma F_m; \quad G'_m = (p_m \gamma - q_m) F_m, \quad (6)$$

where $\gamma = (f_m - p_m)/q_m$.

In order to preserve the accuracy of the functions if it should happen that F_m is close to a zero, i.e. $f_m > \text{ACCUR}^{-1/2}$, then the solutions could be recast to be multiples of F'_m instead of F_m , i.e.

$$\begin{aligned} F &= F' f^{-1}, \quad G = (1 - p f^{-1}) F', \\ G' &= [(1 - p f^{-1}) p/q - q f^{-1}] F', \end{aligned} \quad (7)$$

where F' is determined as

$$F'_m = \text{sign}(F_m) [(1 - p f^{-1})^2/q + q f^{-2}]^{-1/2}. \quad (8)$$

However, provided F_m is not exactly zero in the computer representation (see line 203 of the listing) the method as programmed loses no accuracy in F'_m , G_m and G'_m . As in KLEIN, ref. II, a quantity PACCQ is made available in a COMMON block which provides a guide to the final relative accuracy of the functions (see next section).

The final stage of the calculation is the upward recurrence of G_λ and G' from $\lambda = m$ to $\lambda = l - 1$

means of

$$\begin{aligned} G_{\lambda+1} &= (S_{\lambda+1}G_{\lambda} - G'_{\lambda})/R_{\lambda+1}, \\ G'_{\lambda+1} &= R_{\lambda+1}G_{\lambda} - S_{\lambda+1}G_{\lambda+1}, \end{aligned} \quad (9)$$

(corrected eqs. (26) from I). This is bypassed entirely if $\text{MODE} = 3$ (F_{λ} only required) and only G_{λ} is stored if $\text{MODE} = 2$.

2.2. The subroutine COULFG for Bessel functions

The relationship between the Coulomb functions and the Bessel functions is a very close one, and it is fruitful to think of the Coulomb functions as independent basis functions in that they have a unit Wronskian and remain finite as $x \rightarrow \infty$, behaving in the simplest way as circular functions of unit amplitude, $F_{\lambda} \rightarrow \sin \theta_{\lambda}$ and $G_{\lambda} \rightarrow \cos \theta_{\lambda}$. The angle θ_{λ} is the Coulomb phase

$$\theta_{\lambda} = x - \eta \ln(2x) - \frac{1}{2}\pi\lambda + \arg \Gamma(\lambda + 1 + i\eta).$$

Bessel functions from this view point, I, are merely renormalised Coulomb function for uncharged particles ($\eta = 0$), and the relationships are;

a) *Spherical Bessel functions for real $\lambda > -1$*

$$\begin{aligned} j_{\lambda}(x) &= F_{\lambda}(0, x)/x, \\ y_{\lambda}(x) &= -G_{\lambda}(0, x)/x, \\ j'_{\lambda}(x) &= [F'_{\lambda}(0, x) - x^{-1}F_{\lambda}(0, x)]/x, \\ y'_{\lambda}(x) &= -[G'_{\lambda}(0, x) - x^{-1}G_{\lambda}(0, x)]/x, \text{ and} \end{aligned} \quad (10)$$

b) *Cylindrical Bessel functions for real $\mu > -\frac{1}{2}$*

$$\begin{aligned} J_{\mu}(x) &= (2/\pi x)^{1/2} F_{\mu-1/2}(0, x), \\ Y_{\mu}(x) &= -(2/\pi x)^{1/2} G_{\mu-1/2}(0, x), \\ J'_{\mu}(x) &= (2/\pi x)^{1/2} \\ &\quad \times [F'_{\mu-1/2}(0, x) - \frac{1}{2}x^{-1}F_{\mu-1/2}(0, x)], \\ Y'_{\mu}(x) &= -(2/\pi x)^{1/2} \\ &\quad \times [G'_{\mu-1/2}(0, x) - \frac{1}{2}x^{-1}G_{\mu-1/2}(0, x)]. \end{aligned} \quad (11)$$

They are calculated precisely in this manner in COULFG, controlled by the 'function' parameter

KFN which determines the function; KFN = 0 (Coulomb), = 1 (spherical Bessel), = 2 (cylindrical Bessel). A related paper in this series will describe the companion subroutines BESSJY and SBESJY which evaluate cylindrical and spherical Bessel functions, respectively, and which are progressively more compact, as was outlined in fig. 1 of I.

3. Notes on subroutine COULFG

The calling sequence is:

CALL COULFG(XX, ETA1, XLMIN, XLMAX, FC, GC, FCP, GCP, MODE 1, KFN, IFAIL)

and the named COMMON block contains COMMON/STEED/PACCQ, NFP, NPQ, IEXP, M1. The variables have the following type and meaning (and 'real' indicates REAL*8 on the IBM version). The basic accuracy is set, within 2-3S, by the parameter ACCUR. An appropriate value for the IBM (REAL*8) version is 10^{-16} (56-bit mantissa), and for the CDC (single precision) version it is 10^{-14} (48-bit mantissa); this is set within the code, and although the value could be transmitted as an argument if desired, for most purposes variable accuracy is not needed. For ultra-precise work (REAL*16 variables on an IBM system) then 10^{-33} is a suitable value (112-bit mantissa) while the use of DOUBLE PRECISION variables on CDC machines can allow ACCUR = 10^{-28} (96-bit mantissa).

XX (real) $x > 0$ The accuracy of the Coulomb functions is reduced if $x < x_m$, where x_m is the turning point for the lowest angular momentum, m (eq. (2)). An estimate of the reduction in accuracy, is to be found in the variable PACCQ in the COMMON block.

COULFG fails at once (IFAIL = -1) if $x < \text{ACCUR}^{1/2}$ (= 10^{-8} in the standard version) and no change is made to the output arrays.

ETA1 (real) η Positive, negative or zero. COULFG has been tested for the η -range $10^4 \geq \eta \geq -10^4$. If Bessel functions are requested (KFN = 1,2), then η is set to zero (line 117)

XLMIN (real) m Minimum value of the order (angular momentum) for F_{λ} , j_{λ} or J_{μ} .

XLMAX (real) l Maximum value of the order.

The subroutine fails if $l < m$. The difference between these two quantities must be an integer ≥ 0 . Otherwise the calculation proceeds from $\lambda = m$ to $\lambda = m + \text{INT}|l - m + \text{ACCUR}|$ and an informative message (FORMAT 2040, line 328) is output.

Limitation on m For the function parameter; KFN = 0, 1 then $m > -1$. For cylindrical Bessels, when KFN = 2, then $m > -\frac{1}{2}$.

KFN (integer) Function parameter, i.e.

KFN = 0 Coulomb functions,

KFN = 1 spherical Bessel functions,

KFN = 2 cylindrical Bessel functions.

MODE1 (integer) For each of these choices of KFN there is the option of not calculating or storing the derivatives (MODE1 = 3) or the derivatives and the irregular function (MODE1 = 2). Table 1 contains the details.

FC (N1) These arrays contain the output functions $F_\lambda, G_\lambda, F'_\lambda, G'_\lambda$, arranged by order, and should be dimensioned to at least to the size L1, the integer part of $\text{XLMAX} + 1 = \text{int}(l + 1)$. When MODE1 = 1, all arrays must be at least this size- N1, N2, N3 \geq L1.

N2, N3 \geq 1. The unused arrays are not accessed during the program execution.

IFAIL (integer) Monitors the subroutine execution. A non-zero value indicates a program failure and is accompanied by an error message.

IFAIL = 0 Successful calculation, with no errors detected.

= -1 Input x -value $\text{XX} \leq \text{SQRT}(\text{ACCUR})$. The values are printed according to FORMAT 2000, and control returns to the calling program.

= -2 Problem with the order values XLMIN, XLMAX or XLM (minimum λ for the equivalent Coulomb function), where $\text{XLM} = \text{XLMIN}$ for KFN = 0, 1 and $\text{XLM} = \text{XLMIN} - 0.5\text{DO}$ for KFN = 2.

Either $\text{XLMAX} < \text{XLMIN}$ or $\text{XLM} \leq -1$. The incorrect values are output via FORMAT 2005, and control returns to the calling program.

= 1 The continued fraction CF1 has failed to converge to within a relative accuracy ACCUR, after 20000 iterations (the loop around statement 4). This will occur for $x/(l+1) \gtrsim 20000$. This maximum number can be increased (it is the value of ABORT, line 101) but the subroutine is in an inefficient region [7,8] and asymptotic

Table 1

The choice of the parameters KFN and MODE1 and the functions stored in the arrays FC, GC, FCP, GCP

KFN	MODE1=1	MODE1=2	MODE1=3
0	$F_\lambda(\eta, x), G_\lambda(\eta, x) F'_\lambda(\eta, x) G'_\lambda(\eta, x)$	$F_\lambda(\eta, x), G_\lambda(\eta, x)$	$F_\lambda(\eta, x)$
1	$j_\lambda(x), y_\lambda(x), j'_\lambda(x), y'_\lambda(x)$	$j_\lambda(x), y_\lambda(x)$	$j_\lambda(x)$
2	$J_\mu(x), Y_\mu(x), J'_\mu(x), Y'_\mu(x)$	$J_\mu(x), Y_\mu(x)$	$J_\mu(x)$

Notes

1. The derivatives F', G', j', y', J', Y' , are with respect to the dimensionless variable x .
2. The range of λ for KFN=0, 1, is $-1 < m \leq \lambda \leq l$ and the arrays are filled from the starting element

$\text{M1} = \max(\text{int}(m), 0) + 1$, i.e. $\text{M1} \geq 1$.

For KFN=2 then the range of the order μ is

$$-\frac{1}{2} < m \leq \mu \leq l,$$

starting from M1 as defined above. The value of M1 is available in the COMMON block.

3. The internal parameter XLM, lines 126–127, is, in each case, the equivalent Coulomb function angular momentum (eqs. (10), (11)).

methods, eg. ref. [4], might well be preferable. Values of the relevant variables are output with FORMAT 2010, e.g. the maximum λ -value of XLM + int($l-1 + \text{ACCUR}$) and control returns to the calling program with the arrays undefined.

- = 2 The continued fraction CF2 has failed to converge after 20000 iterations (the loop around statement 8), according to the criterion (line 242).

$$|dp| + |dq| \leq \text{ACCUR} \times (|p| + |q|).$$

The failure may occur for $x \ll x_m$ and values of the relevant variables are output according to FORMAT 2020. Control returns to the calling program with arrays FC and possibly, FCP and GC modified, and GCP not defined.

- = 3 Apparent convergence of CF2 but either

- a) q is negative, or
- b) q is so small ($< 10^4 \times \text{ACCUR} \times |p|$, i.e. $10^{-10} \times |p|$) that its value and the 'convergence' are suspect. Control returns as for IFAIL = 2.

COMMON/STEED/PACCQ, NFP, NPQ, IEXP,

M1 contains the following variables:

PACCQ (real) has the value:

$$\begin{aligned} & \frac{1}{2} \text{ACCUR}/q \text{ if } |p/q| \leq 1, \\ & |p| \times \frac{1}{2} \text{ACCUR}/q \text{ if } |p/q| > 1, \\ & \frac{1}{2} \text{ACCUR} \text{ if } q > 1. \end{aligned}$$

It is a rough guide to the final relative accuracy and hence monitors the number of significant digits lost as the magnitude of q falls greatly below that of p [7,8]. If $G_m > (10^4 \text{ACCUR})^{-1/2}$ then PACCQ = 1.0DO, and JWKB approximations are used.

NFP (integer) The number of iterations required to evaluate CF1.

NPQ (integer) The number of iterations required to evaluate CF2.

These two qualities are a crude measure of the efficiency of the code for the given input parameters [8].

IEXP (integer) If IEXP = 1 then JWKB was not used and $x(2\eta - x) > \max(m^2 + m, 0)$. If IEXP = 0 and $G > 10^6$, or if IEXP > 70, then the

JWKB estimates for G_m and F_m are used and should be scaled by $10^{\pm \text{IEXP}}$. The constant, 70, is close to the underflow exponent on the IBM compiler.

M1 (integer) The starting array element for the function arrays. Thus if the Coulomb functions for $m=350$ to $l=449$ were required in the course of a larger calculation (see e.g. ref. [6]) then M1 = 351 and L1 = M1 + 99 = 450 (see line 138).

3.1. Informative message

During the calculation of CF1 it is possible for a loss of significant digits to occur at the step (line 168)

$$D = \text{TK} - D * (\text{ONE} + \text{EK} * \text{EK}).$$

i.e. $D_k^{-1} = T_k - D_{k-1}(1 + \eta^2/k^2)$, and information that this has happened is printed out (FORMAT 1000) together with appropriate variables. It does not follow that the loss will propagate through to the final f -value [7], or to the Coulomb functions. However, after a second such occurrence, the subroutine fails with IFAIL = 1 since for an extremely large value of x a semi-infinite loop can be encountered. Further remarks are given in II, section 4.1. An additional complication with COULFG arises with the recurrence between the maximum and minimum λ -values.

3.2. Range errors during recurrence in order

During the downward recurrence, DO LOOP 6, for F_λ, F'_λ it is possible for overflows to occur if the λ -range is large enough. An offset of $10^{-30} = \text{TM30}$ is provided so that a range of $\approx 10^{105}$ is available on the IBM compiler, as programmed. However, upon renormalization by W (lines 296, 297) underflows are unavoidable. Similarly during the upward recurrence for G_λ, G'_λ , in DO LOOP 12, overflows are inevitable after the value of G'_λ reaches $\approx 7.2 \times 10^{75}$ (the limit for the IBM compiler). On the CDC machine the limits are raised to $\approx 10^{290}$ although eventually the same errors will occur. The overflow and underflow errors are not trapped in COULFG on the grounds that normal usage will not require such small values of x/x_λ .

The problems of overflows during the recurrence relations can be overcome by the addition of a scaling factor, for example in powers of $10^{\pm 60}$, which could be extracted from the Coulomb functions as soon as G'_λ , exceeded a power of 10^{60} or F'_λ , F'_λ , was less than a power of 10^{-60} .

4. Test calculations with subroutine COULFG

4.1. Test output for COULFG

The test sequence for COULFG examines the use of the subroutine for each function, for a wide range of arguments and for most of the error conditions. The tests were run on an IBM 370/165 system at Daresbury Laboratory and on the replacement NAS 7000 system, using both the G and the HX compilers and REAL*8 variables; they were also verified on a CDC 7600 system at the University of Manchester Regional Computing Centre using single-precision variables and on a GEC 4065 machine.

Examples which appear in Bardin et al. [4] are given for a range of λ -values, for both positive and negative values of η ; some tests from Kölbig's review [9], a few calculations from earlier papers [12–14], and an example from the heavy-ion DWBA program A-THREE of Auerbach [15] using parameters appropriate to the scattering of ^{18}O on ^{64}Ni at 65 MeV.

For each set (range of λ) the values of NFP and NPQ are given, the first being the number of iterations to evaluate CF1 at $\lambda = l$ (the maximum value) and the second referring to CF2 at $\lambda = m$ (the minimum value). If the value of PACCQ is greater than $\frac{1}{2}\text{ACCUR}$ ($= 5 \times 10^{-17}$ in this case), then the value of x is less than the turning point for the lowest partial wave, x_m , and accuracy begins to be lost for this reason. The value of PACCQ is very approximately the relative final accuracy of the Coulomb wavefunctions [1,7,8]. An example is for $\eta = 10$, $x = 5$, $\lambda = 0-20$ when $\text{PACCQ} = 2 \times 10^{-6}$.

In cases when the x -value is less than x_m , the JWKB approximation will be used if $G_m > (100 \text{ ACCUR}^{1/2})^{-1}$. These cases are identified by $\text{PACCQ} = 1.0$, $\text{NPQ} = 0$ and by a value of the

exponent IEXP which is either 0 or > 70 (on the IBM system). The user should check the value of IEXP since the arrays F_λ , F'_λ should be multiplied by $10^{-\text{IEXP}}$ and the arrays G_λ , G'_λ by $10^{+\text{IEXP}}$. The rest of the calculation continues in the normal way using, where necessary, the upward recurrence relations (9).

The use of the MODE parameter is not illustrated but it can be readily verified by setting the function arrays to zero before the call to COULFG, and then using, in turn, $\text{MODE} = 2, 3$ (section 3). There is no change in the derivative arrays when $\text{MODE} = 2, 3$ or in the irregular function array G_λ when $\text{MODE} = 3$.

The Bessel-function tests, which follow in the output, are illustrative only as a full discussion is deferred until the next paper of this series, which will describe the subroutine BESSJY and its MODE-3 version BESSJ. Some of the standard tabulations in Abramowitz and Stegun are recalculated for the cylindrical Bessel [16] and the spherical Bessel [17] cases. The well-known difficulties [1] experienced by non-direct methods in obtaining the irregular function $Y_\mu(x)$ or $y_\lambda(x)$ are emphasised by noting that there are errors in the last two or three places on occasion in refs. [16,17]. Bessel functions of order $1/3$, $J_{1/3}$, $Y_{1/3}$ are compared with the 7D values of table III in Watson [18] and complete agreement is found. Once again the trends of NFP and NPQ emphasise the general conclusions drawn in refs. [1,8] that the singularity in the irregular function at $x = 0$ forces a rapid increase of NPQ as $x \rightarrow 0$ and that NFP increases approximately as x does. The general methods of COULFG are most effective in the "non-asymptotic" region where x is neither too small nor too large. The JWKB approximation for Bessel functions is contained within the Coulomb formula [2,8] (by setting $\eta = 0$) and is invoked in cases where both $G_m > (100 \text{ ACCUR}^{1/2})^{-1}$ and $x^2 < m(m+1)$ for J_m , y_m , and $x^2 < m^2 - \frac{1}{4}$ for J_m , Y_m , but not for J_0 and Y_0 .

The use of the recurrence relations for a λ -value where one of the functions has an approximate zero does not increase the errors for any of the functions: a test example of Blanch [19] is a convenient example of the stability of the methods of COULFG.

Table 3
Curtis calculations

X	PL(ETA, X)	Z	PL(A, Z)	P'L(A, Z)	QL(A, Z)	Q'L(A, Z)	NFP	NPQ
A-VALUE = 2.000 ETA = -0.707107 L = 0								
1.414	0.424398	1.0	0.141533	-0.465728	0.247936	0.308650	8	60
2.828	-0.872102	2.0	-0.290838	-0.152161	0.082660	-0.503983	12	31
4.243	-0.160995	3.0	-0.053690	0.497580	-0.309635	-0.094746	14	22
5.657	0.936694	4.0	0.312378	0.074412	-0.045234	0.498719	17	17
7.071	0.142960	5.0	0.047676	-0.487531	0.318892	0.077302	19	14
8.485	-0.946141	6.0	-0.315529	-0.091492	0.059404	-0.487182	22	12
9.899	-0.226086	7.0	-0.075398	0.473737	-0.317380	-0.116720	24	11
11.314	0.927308	8.0	0.309248	0.145221	-0.097265	0.468976	26	10
12.728	0.358766	9.0	0.119645	-0.449815	0.305600	0.181299	28	9
14.142	-0.874426	10.0	-0.291613	-0.215643	0.146682	-0.437306	30	8
A-VALUE = 2.000 ETA = -0.707107 L = 1								
1.414	0.806284	1.0	0.465728	0.424598	-0.308650	0.743809		
2.828	0.515181	2.0	0.297580	-0.723723	0.462653	0.479306		
4.243	-0.799461	3.0	-0.461786	-0.468929	0.301169	-0.728125		
5.657	-0.534425	4.0	-0.308696	0.705613	-0.464793	-0.484297		
7.071	0.777999	5.0	0.449390	0.502540	-0.332416	0.690743		
8.485	0.613607	6.0	0.354433	-0.651225	0.437678	0.542944		
9.899	-0.708266	7.0	-0.409110	-0.576859	0.388760	-0.618917		
11.314	-0.719871	8.0	-0.415813	0.563908	-0.383869	-0.627681		
12.728	0.594617	9.0	0.343464	0.664236	-0.452943	0.514183		
14.142	0.827694	10.0	0.478095	-0.444553	0.305292	0.714808		
A-VALUE = 2.000 ETA = -0.707107 L = 2								
1.414	0.316231	1.0	0.608874	1.415327	-2.681744	2.479364		
2.828	1.007017	2.0	1.938918	0.518441	-0.551065	2.588800		
4.243	0.452389	3.0	0.871032	-2.454104	1.729599	1.217576		
5.657	-0.814388	4.0	-1.568029	-1.543479	1.076216	-2.323967		
7.071	-0.631876	5.0	-1.216619	2.125287	-1.461114	-1.808190		
8.485	0.683438	6.0	1.315896	1.991482	-1.368646	1.960285		
9.899	0.766966	7.0	1.476723	-1.729111	1.190247	2.198852		
11.314	-0.530860	8.0	-1.022122	-2.334597	1.608107	-1.517316		
12.728	-0.876747	9.0	-1.688096	1.248405	-0.863035	-2.504448		
14.142	0.347545	10.0	0.669166	2.591224	-1.791991	0.988865		
A-VALUE = 0.200 ETA = -2.236068 L = 0								
0.447	0.447884	1.0	0.267189	-0.230547	0.200725	0.422468	6	180
0.894	-0.186654	2.0	-0.111350	-0.406813	0.368443	-0.083227	7	94
1.342	-0.668542	3.0	-0.398824	-0.133387	0.118312	-0.359491	9	65
1.789	-0.613009	4.0	-0.365695	0.184173	-0.241068	-0.313804	10	50
2.236	-0.144028	5.0	-0.085921	0.340922	-0.446516	-0.080629	11	40
2.683	0.414747	6.0	0.247421	0.294976	-0.397197	0.169717	12	35
3.130	0.765813	7.0	0.456851	0.108862	-0.143252	0.314239	14	30
3.578	0.760264	8.0	0.453541	-0.112429	0.179676	0.306376	14	27
4.025	0.423852	9.0	0.252852	-0.272584	0.426167	0.170015	15	24
4.472	-0.087295	10.0	-0.052077	-0.315866	0.499464	-0.026716	16	22
A-VALUE = 0.200 ETA = -2.236068 L = 1								
0.447	0.352790	1.0	0.576367	0.801566	-1.056169	0.602174		
0.894	0.707714	2.0	1.156220	0.244060	-0.252486	0.979087		
1.342	0.610975	3.0	0.998174	-0.531021	0.701541	0.822631		
1.789	0.137871	4.0	0.225246	-0.928151	1.236514	0.204181		
2.236	-0.416507	5.0	-0.680464	-0.802134	1.094605	-0.463863		
2.683	-0.766890	6.0	-1.252898	-0.301820	0.403201	-0.855591		
3.130	-0.765802	7.0	-1.251121	0.298165	-0.478628	-0.840008		
3.578	-0.435227	8.0	-0.711049	0.738456	-1.158981	-0.475080		
4.025	0.073185	9.0	0.119565	0.864837	-1.372077	0.058878		
4.472	0.555068	10.0	0.906838	0.659924	-1.057003	0.547088		
A-VALUE = 0.200 ETA = -2.236068 L = 2								
0.447	0.057470	1.0	0.699823	1.832099	-24.293641	31.159617		
0.894	0.304806	2.0	3.711672	3.925266	-12.281445	4.878293		
1.342	0.636330	3.0	7.748707	3.699421	-7.841190	4.814571		
1.789	0.846894	4.0	10.312784	1.126228	-2.268677	6.182570		
2.236	0.793999	5.0	9.668670	-2.435452	3.937810	5.866806		
2.683	0.465932	6.0	5.673740	-5.318868	8.759894	3.475985		
3.130	-0.027435	7.0	-0.334079	-6.327195	10.473014	-0.148924		
3.578	-0.511608	8.0	-6.229937	-5.112727	8.498059	-3.670391		
4.025	-0.819428	9.0	-9.978323	-2.173931	3.580603	-5.865773		
4.472	-0.850384	10.0	-10.355285	1.427603	-2.555408	-6.051639		

4.2. High-accuracy test calculations

Test calculations using the high-accuracy mode of COULFG, with ACCUR = 10^{-33} , are given in table 2 and these represent the current convenient accuracy limit of the code. Table 2 contains examples of Bessel functions of order $\mu = 1/3$ taken from the exceptional work of Makinouchi, who lists in ref. [20] values of $Y_{1/3}(x)$ between $x = 0.01$ and $x = 100$ to 30 decimals. It should be noted that these were computed from his values of $J_{1/3}(x)$ and are not obtained directly as in the present method. Nevertheless the agreement is virtually complete with the maximum discrepancy being in the last digit. At the same time, of course, COULFG also produces $J_{1/3}(x)$, $J'_{1/3}(x)$ and $Y'_{1/3}(x)$ to the same accuracy. The calculation of Airy functions using COULFG (or KLEIN) and Bessel functions of order $1/3$ is discussed in I and also in a subsequent paper of this series.

The equally remarkable calculations of Aldis in 1900 [21], who obtained $J_0(x)$, $J_1(x)$, $-\frac{1}{2}\pi Y_0(x)$, $-\frac{1}{2}\pi Y_1(x)$ to 21D, for $x = 0.1$ (0.1) 6.0 are illustrated in table 2. A few of his x -values are displayed and every digit is correct for his complete table. The irregular function is that defined by Heine (ref. [18] p. 65) and is of no particular current relevance.

Finally, a few of the high-precision results of Strecok and Gregory [21] for $\lambda = 0$ are given; here values for $\lambda = 10$ are also quoted. The full comparison was made in ref. [7].

4.3. The results of Curtis for negative η

Table 3 is a comparison of some of the results of Curtis [23] which are appropriate to electron scattering at positive energies. The relation between the functions of Curtis, $P_L(a, z)$, $Q_L(a, z)$ and the Coulomb functions $F_\lambda(\eta, x)$, $G_\lambda(\eta, x)$ is given by his equations (2.74)–(2.77) and (6.15)–(6.16) namely;

$$\begin{aligned} P_L(a, z) &= K_L(a) F_L(\eta, x), \\ Q_L(a, z) &= K_L(a)(1 - e^{2\pi\eta})^{-1} G_L(\eta, x), \end{aligned} \quad (12)$$

where

$$\eta = -a^{-1/2}, \quad x = za^{1/2}, \quad z = -\eta x,$$

$$\begin{aligned} K_0 &= a^{-1/4}(2\pi)^{-1/2}(1 - e^{2\pi\eta})^{1/2}, \\ K_1 &= 3(1 + a)^{-1/2} K_0, \\ K_2 &= 10(1 + 4a)^{-1/2} K_1. \end{aligned} \quad (13)$$

In the table results are given for $a = +2.0$ ($\eta = -0.7071$) and $a = +0.2$ ($\eta = -2.2361$) for a range of z -values and for both forms of the Coulomb functions. The agreement between the calculations is exact.

5. Comments on the program

5.1. Modifications to COULFG

The ultimate accuracy of Steed's method, around which COULFG is based, is not that of table 2 (i.e. $\approx 10^{-33}$) but is limited by the effective word-length of the computer used. In principle, then, arbitrary accuracy can be achieved with a suitable multiple-precision package, such as that of Schonfelder [24] or of Brent [25], and these could be exploited to provide an enhancement of COULFG or of KLEIN for research purposes. Indeed, such efforts are being made by the author for the Γ -function subroutine [1].

In normal use when the value of PACCQ, which approximates the resultant relative accuracy, becomes too large for the application in hand, i.e. when x becomes significantly less than x_m , then one has recourse to various alternatives:

1) the integration which was removed from RCWFN [3], lines 60–67, 107 and 149–182, could be restored;

2) the value of ACCUR could be decreased with the use of extended precision (at the cost of significantly greater computing time);

3) the asymptotic methods exploited by Bardin et al. [4] could be used; or

4) the JWKB estimates employed in KLEIN [2] could be adopted. This last alternative is programmed in COULFG and the resultant accuracies are in general better than 1%.

To retain high accuracy over an extended integration range is not really possible even with the 4th-order Runge–Kutta method [3], and improved techniques such as those of Strecok and Gregory [22] are required.

5.2. Integer values of the order

Since calculations with λ integral are likely to be a frequent requirement it may be convenient to detail here the modifications necessary to convert the subroutine to accept integer values of the Coulomb or Bessel order parameter. They are given in table 4.

5.3. Modifications for complex-variable programming

In operating the code on a machine with adequate single-precision word-length (e.g. a CDC

Table 4
Possible modifications to COULFG

```

***          INTEGER VALUES OF THE ORDER          ***
CHANGE XLMIN,XLMAX TO LMIN,LMAX AT LINES 55,126,128,137,309,310
CHANGE FORMAT IN LINE 311 TO 2110,1FD15.6
DELETE LINES 71 - 77,131,132,328
REPLACE LINE 70 BY 'AND INTEGER LAMBDA .GE. 0'
CHANGE LINE 133 TO LXTRA = LMAX - LMIN
CHANGE LINE 138 TO M1 = LMIN + 1
IN CALLING PROGRAM LINE 30 MAKE THE LAMBDA VALUES INTEGRAL BY, FOR
EXAMPLE, CHANGING XM,XL(JMAX) INTO IDINT(XM),IDINT(XL(JMAX))

*** COMPLEX-VARIABLE PROGRAMMING WITH SINGLE-PRECISION VARIABLES ***
TO CONVERT COULFG INTO SINGLE PRECISION, DELETE 95,332, CHANGE 'D' INTO
'E' EXPONENTS IN LINES 101 - 103,107 & FUNCTION NAMES LISTED IN 100,335
DELETE DZERO & 0.0D0 IN 337 AND CHANGE DZERO TO ZERO IN 341
ADD FOLLOWING DECLARATION AFTER LINE 97
COMPLEX C1,AA,BB,DD,DL,PQ,TWOI
INSERT NEXT TWO LINES AFTER LINE 106 - DATA & LOCAL ARITHMETIC FUNCTION
DATA C1,TWOI / (1.0E0,0.0E0) , (0.0E0,2.0E0) /
CXMOD(PQ) = ABS(REAL(PQ)) + ABS(AIMAG(PQ))
NOTE THAT SINGLE PRECISION CONSTANTS AND FUNCTIONS ARE USED

REPLACE LINES 215 - 242 BY THE NEXT 15 LINES
PQ = CMPLX(ZERO,ONE - ETA*XI)
AA = CMPLX(-EZMMI,ETA)
BB = CMPLX(TWO*(X - ETA),TWO)
DD = C1/BB
DL = AA*DD*CMPLX(ZERO,XI)
8 PQ = PQ + DL
PK = PK + TWO
AA = AA + CMPLX(PK,M1)
BB = BB + TWOI
DD = C1/(AA*DD + BB)
DL = DL*(BB*DD - C1)
IF(PK .GT. TA) GO TO 120
IF(CXMOD(DL) .GE. CXMOD(PQ)*ACC) GO TO 8
P = REAL(PQ + DL)
Q = AIMAG(PQ + DL)

REPLACE OUTPUT LINES 319,324 BY NEXT TWO LINES
WRITE (6,2020) ABORT,PQ,DL,ACC
WRITE (6,2030) PQ,ACC,LXTRA,M1

***          NON-ANSI DOUBLE-PRECISION COMPLEX VARIABLES          ***

CHANGE FUNCTION TYPES ABOVE COMPLEX , CMPLX, ABS, REAL,AIMAG
BECOME COMPLEX*16,DCMPLX,DABS,DREAL,DIMAG

NOTE THAT EXTENDED-PRECISION WITH COMPLEX*32 VARIABLES IS AVAILABLE
ON THE IBM HX COMPILER SIMPLY BY THE USE OF THE AUTODOUBLE FACILITY
A SUITABLE ACCUR IS THEN 10**(-33) FOR THE 112-BIT MANTISSA
MAKE SURE THE EXPONENTS OF LINES 102, 103 READ 'Q0'

*** IF RESULTS OF DOUBLE-PRECISION CALCULATIONS ARE TO BE STORED IN
SINGLE-PRECISION ARRAYS ***

REPLACE LINE 288 WHICH EXTRACTS RL FROM STORAGE(NOW SINGLE PRECISION) BY
IF(ETANEQ) RL = DSQRT(ONE + EL*EL)

```

computer with a 48-bit mantissa), the evaluation of CF2 can be replaced by rather more compact and transparent coding as is detailed in table 4. Use is made of a local arithmetic function CXMOD which does not use the square-root function. Certain compilers support NON-ANSI DCMPLX variables, i.e. COMPLEX*16, and indeed the IBM HX compiler allows extended-precision complex functions. Suitable changes for these are also given in the table.

5.4. Single-precision arrays

It occasionally happens that restricting the Coulomb function arrays to REAL*4 accuracy is thought to be adequate (though probably not for matching to internal wavefunctions) while their calculation should be carried out in REAL*8 variables. For such a situation the only change that needs to be made to COULFG is to avoid the use of the stored value $RL = (1 + \eta^2/\lambda^2)^{1/2}$, which is in the array GC, since it is there truncated to REAL*4. Instead the value should be recalculated at the point where it is required in the upward recurrence; that is, replace in line 288 $RL = GC(L + 1)$ by $RL = DSQRT(ONE + EL * EL)$. This modification will preserve REAL*8 accuracy of the irregular functions until the final truncation by storage. In this respect the coding in COULFG is superior to that in RCWFN.

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- [11] A number of important errors in ref. [1] should be corrected:
- a) the quantity F_{LMAX} after eq. (25), before eq. (34), and in section 5 should be read as $\text{sign}(F_m)$;
 - b) in the upward recurrence relations eq. (26), (36), for the quantity S_λ read $S_{\lambda+1}$;
 - c) the angular momentum term in square brackets in eq. (46) is indeed negative.
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TEST RUN OUTPUT

TEST OF THE MANCHESTER COULOMB WAVEFUNCTION PROGRAM - COULFG

ETA	X	XL	F	G	FP	GP	NFP	NPQ	PACQ
-50.000	5.000	0.0	1.5223697514D-01	4.41680690236D-01	2.03091041166D+00	-6.76485374767D-01			
-50.000	5.000	10.0	-3.68114360218D-01	3.31588324611D-01	1.33846751032D+00	1.51088862814D+00			
-50.000	5.000	50.0	4.20266913301D-23	1.32186136031D+21	3.83028028856D-22	-1.17470596213D+22	7	45	5.D-17
-50.000	50.000	10.0	2.77129451057D-01	-7.10519441453D-01	-1.22086610371D+00	-4.78761065984D-01			
-50.000	50.000	20.0	7.55876294797D-01	-1.51281009566D-01	-2.52360602589D-01	-1.27246037462D+00	91	14	5.D-17
-4.000	5.000	0.0	4.07862723006D-01	6.74327035383D-01	1.09821233636D+00	-6.36110427280D-01			
-4.000	5.000	5.0	2.83657834139D-01	-8.70189296614D-01	-1.04544995379D+00	-3.18202528554D-01			
-4.000	5.000	10.0	7.21985521560D-02	5.28003488054D+00	1.10084353023D-01	-5.79998855560D+00	11	21	5.D-17
-4.000	50.000	50.0	1.59533986051D+00	-5.15310138675D-03	-8.64299712691D-02	-6.26546507323D-01			
-4.000	50.000	100.0	3.14179162347D-19	9.37788817811D+17	5.37489508885D-19	-1.57855264866D+18	13	28	1.D-16
10.000	5.000	0.0	1.72074051358D-06	1.67638043033D+05	3.09759065782D-06	-2.79371561379D+05			
10.000	5.000	5.0	1.31825659927D-07	1.84916325924D+06	2.80452562756D-07	-3.65177329860D+06			
10.000	5.000	20.0	4.04282716774D-18	2.77871505109D+16	1.83778234798D-17	-1.21037242653D+17	8	61	2.D-06
10.000	50.000	30.0	-3.52635436430D-01	-1.39731750174D+00	-6.64436751932D-01	2.02960027068D-01			
10.000	50.000	50.0	5.76729101127D-03	1.34226818062D+02	3.89565495430D-03	-8.27249102369D+01	22	20	1.D-16
100.000	1000.000	0.0	-1.65581311975D-01	-1.04432596294D+00	-9.34063273097D-01	1.48165711330D-01			
100.000	1000.000	50.0	-1.02728883766D+00	2.53963904213D-01	2.26855983270D-01	9.17353264488D-01			
100.000	1000.000	100.0	9.36270607951D-01	-4.98554253164D-01	-4.43161964632D-01	-8.32088192320D-01	868	6	6.D-17

JWKX EXAMPLES CALCULATIONS OF BARDIN ET AL.

100.000	1.000	0.0	9.01215659555D-03	3.93122386932D+00	1.29458981023D-01	-5.44897458484D+01			
100.000	1.000	5.0	1.07533483439D-03	0.7107311677D+01	1.6586463122D-02	-4.56246324475D+02			
100.000	1.000	10.0	4.72013782540D-06	5.02336440657D+03	8.46148386724D-05	-1.03881286724D+05			
100.000	1.000	30.0	7.10143042253D-26	2.09518618848D+23	2.41881682995D-24	-6.94526405033D+24	5	0	1.D+00
**** IEXP = 123 F,FP *10**(-IEXP) G,GP *10**(IEXP)									
100.000	5.000	0.0	1.63540865026D-02	4.89523124695D+00	1.02979707085D-01	-3.03221167381D+01			
100.000	5.000	20.0	4.77275504854D-08	1.40230439512D+06	3.59703975637D-07	-1.03836364737D+07	12	0	1.D+00
**** IEXP = 108 F,FP *10**(-IEXP) G,GP *10**(IEXP)									
100.000	50.000	0.0	1.70424082262D-54	1.69384477356D+53	2.96321541644D-54	-2.92257584011D+53			
100.000	50.000	20.0	4.53710831483D-56	6.19148506714D+54	8.10665703799D-56	-1.09778631121D+55	31	0	1.D+00
10.000	5.000	30.0	3.2073524701D-28	2.45851379719D+26	2.06510422204D-27	-1.53488133636D+27			
10.000	5.000	50.0	2.15139724500D-52	2.26795313756D+50	2.22594022701D-51	-2.30160742729D+51	6	0	1.D+00
-4.000	5.000	30.0	4.63196983554D-19	1.83494540622D+17	2.77366206877D-18	-1.06012813184D+18			
-4.000	5.000	50.0	4.43423440067D-43	1.13093510555D+41	4.46626563407D-42	-1.11607643453D+42	6	0	1.D+00
500.000	0.001	0.0	6.15007019274D-05	7.51159667969D+00	8.81383260404D-02	-5.49490383314D+03	5	0	1.D+00
**** IEXP = 679 F,FP *10**(-IEXP) G,GP *10**(IEXP)									
500.000	0.025	0.0	4.87216122025D-04	5.11375617981D+00	1.02722117830D-01	-9.74319431732D+02	10	0	1.D+00
**** IEXP = 676 F,FP *10**(-IEXP) G,GP *10**(IEXP)									

OTHER TEST EXAMPLES OF COULOMB FUNCTIONS

2.000	10.000	0.0	-1.06160862047D+00	-3.99306648097D-01	-2.93530453784D-01	8.31560069659D-01			
2.000	10.000	1.0	-8.65737387536D-01	-7.46893228412D-01	-5.55780527035D-01	6.75598970652D-01	23	15	6.D-17
4.000	10.000	1.0	1.37343785651D+00	5.2328665858D-01	1.37992480244D-01	-6.75519882474D-01			
4.000	10.000	4.0	8.89699380561D-01	1.56161024485D+00	3.43085462667D-01	-5.21787737277D-01	18	21	1.D-16
5.000	10.000	0.0	9.17944918946D-01	1.60852455560D+00	3.31032101931D-01	-5.09318942458D-01			
5.000	10.000	1.0	8.53200693407D-01	1.70871952005D+00	3.29295517733D-01	-5.12571455183D-01			
5.000	10.000	7.0	8.94776433905D-02	7.63318878690D+00	7.46133154652D-02	-4.81083833599D+00	16	26	2.D-16
10.000	10.000	0.0	1.62627112503D-03	3.07873216608D+02	1.70604763209D-03	-2.91927723806D+02			
10.000	10.000	7.0	9.92870627105D-05	4.03294624069D+03	1.28959191579D-04	-4.83360570871D+03			
10.000	10.000	10.0	8.18312960548D-06	4.21630275894D+04	1.22688004233D-05	-5.89894924245D+04	14	40	4.D-12
25.000	10.000	0.0	1.54828713210D-16	1.61423768840D+15	1.61423768840D+15	-3.17484176641D+15			
25.000	10.000	1.0	1.427671057658D-16	1.74648170081D+15	2.90760570837D-16	-3.44772363283D+15	21	0	1.D+00
5.000	5.000	0.0	2.76730116686D-02	1.81934951590D+01	3.03600120489D-02	-1.61762396201D+01			
5.000	5.000	1.0	2.22669570738D-02	2.17262189626D+01	2.53170497141D-02	-2.02073517694D+01			
5.000	5.000	15.0	1.56591550129D-13	7.93097863019D+11	6.46100173074D-13	-3.11370142920D+12			
5.000	5.000	20.0	1.88142640121D-14	6.29076034991D+12	8.13652532644D-14	-2.59182249219D+13	8	45	1.D-14
4.000	6.000	0.0	2.94449215620D-01	3.01378381329D+00	2.09173395936D-01	-1.25521341389D+00	17	35	2.D-16
4.000	7.000	0.0	5.51582922742D-01	2.11648513364D+00	3.01934534735D-01	-6.54407580471D-01	18	31	6.D-17
4.000	8.000	0.0	8.24526544489D-01	1.55260815154D+00	3.46666242351D-01	-5.23274542380D-01	19	27	2.D-16
1.000	0.600	0.0	1.07096894447D-01	2.79244038685D+00	2.51925443598D-01	-2.76864439770D+00	7	166	4.D-16
1.000	3.000	0.0	1.08405268420D+00	6.27039514889D-01	3.01916785821D-01	-7.47829194001D-01	13	38	8.D-17
1.000	6.000	0.0	-1.67183367846D-01	-1.09083851821D+00	-8.92504992651D-01	1.58037110148D-01	19	20	6.D-17
1.995	1.200	0.0	4.34422920587D-02	7.46482788802D+00	8.07579506900D-02	-9.13545418683D+00	9	103	3.D-15
3.981	2.400	0.0	5.62399799825D-03	5.82540355738D+01	9.43642344138D-03	-8.00658629852D+01	11	70	2.D-13
18.564	116.790	0.0	-2.19713365418D-01	-1.07820215706D+00	-8.90272795108D-01	1.82537606990D-01			
18.564	116.790	30.0	1.06071645759D+00	3.88736687223D-01	3.02917093766D-01	-8.31744436652D-01			
18.564	116.790	60.0	-1.23263404627D+00	-1.86707561606D-01	-1.14689834856D-01	7.93898678648D-01			
18.564	116.790	80.0	-8.25210276064D-01	1.22923127192D+00	5.71431018640D-01	-3.60641083568D-01			
18.564	116.790	100.0	5.02651913712D-01	4.21053181928D+00	1.44949393217D-01	-7.75260089628D-01			
18.564	116.790	120.0	6.10641818529D-05	1.32512621382D+04	3.8570932195D-16	-8.00586627116D+03			
18.564	116.790	140.0	5.59982983995D-11	1.02087306246D+10	4.94812366777D-11	-8.8688612460D+09			
18.564	116.790	160.0	3.08072086424D-18	3.08072086424D+17	1.62844830136D-18	-3.3727811779D+17			
18.564	116.790	180.0	2.1472540027D-27	1.78257094264D+26	2.81866160498D-27	-2.31719302741D+26			
18.564	116.790	200.0	2.64572896764D-37	1.25560168697D+36	3.99743384864D-37	-1.88246626253D+36	15	8	6.D-17
-4.000	5.000	0.0	4.07862723006D-01	6.74327035383D-01	1.09821233636D+00	-6.36110427280D-01			
-4.000	5.000	1.0	-6.42256329143D-01	-4.67204210149D-01	-7.58912963040D-01	1.00494559453D+00			
-4.000	5.000	2.0	7.98957414374D-01	-1.15121213166D-01	-1.57796947944D-01	-1.22889431433D+00			
-4.000	5.000	3.0	-2.56863093558D-01	7.87989922393D-01	1.14322942201D+00	3.85990587811D-01	17	21	5.D-17

CYLINDRICAL BESSEL FUNCTIONS *** ABRAMOWITZ & STEGUN PAGES 407 - 408

0.0	1.000	0.0	7.65197686558D-01	8.82569642157D-02	-4.40050585745D-01	7.81212821300D-01			
0.0	1.000	10.0	2.63061512369D-10	-1.21618014279D+08	2.61863505622D-09	1.20939993785D+09			
0.0	1.000	30.0	3.48286979425D-42	-3.04812878323D+39	1.04429904344D+40	9.13912933615D+40			
0.0	1.000	40.0	1.10791585113D-60	-7.18484749680D+57	4.43031209141D-59	2.87302862548D+59	4	85	5.D-17
0.0	2.000	0.0	2.23890779141D-01	5.10375672650D-01	-5.76724807757D-01	1.07032431541D-01			

0.0	2.000	10.0	2.51538628272D-07	-1.29184542208D+05	1.23465029377D-06	6.31362881664D+05			
0.0	2.000	30.0	3.65025626647D-33	-2.91322384822D+30	5.46359748625D-32	4.3597776972D+31			
0.0	2.000	50.0	3.22409583944D-65	-1.97615057652D+62	8.05391545651D-64	4.93634176475D+63	4	45	5.D-17
0.0	5.000	0.0	-1.77596771314E-01	-3.08517625249D-01	3.27579137591D-01	-1.47863143391D-01			
0.0	5.000	10.0	1.46780264731D-03	-2.51291100956D+01	2.58466778448D-03	4.24943370028D+01			
0.0	5.000	30.0	2.67117727825D-21	-4.028956841855D+18	1.58102720945D-20	2.382140344474D+19			
0.0	5.000	50.0	2.29424761595D-45	-2.78883701758D+42	2.28297467652D-44	2.77457024037D+43	6	19	5.D-17
0.0	10.000	0.0	-2.45935764451D-01	5.56711672836D-02	-4.34727461689D-02	-2.49015424207D-01			
0.0	10.000	10.0	2.07486106633D-01	-3.59814152183D-01	8.43695786318D-02	1.60514886378D-01			
0.0	10.000	30.0	1.55109607826D-12	-7.25614231610D+04	4.39647875200D-12	2.04761666074D+10			
0.0	10.000	50.0	1.58123276693D-02	-3.64106650180D+27	8.74593525529D-30	1.78297576232D+28	7	10	5.D-17
0.0	50.000	0.0	5.58123276693D-02	-9.80649954701D-02	9.75118281252D-02	5.67956685620D-02			
0.0	50.000	10.0	-1.13847849149D-01	5.72389718205D-03	-4.42289121408D-03	-1.11614574783D-01			
0.0	50.000	30.0	4.84342572455D-02	-1.16457234935D-01	9.24533752861D-02	4.05814213506D-02			
0.0	50.000	100.0	1.11592736908D-21	-3.29380018820D+18	1.93650320925D-21	5.69386591665D+18	13	4	5.D-17
0.0	100.000	0.0	1.99858503042D-02	-7.72443133651D-02	7.71453520141D-02	2.03723120028D-02			
0.0	100.000	10.0	-5.47321769355D-02	5.83315742364D-02	-5.77635437125D-02	-5.47531535058D-02			
0.0	100.000	30.0	8.14601295812D-02	6.13883921201D-03	-6.30376589012D-03	7.76760355155D-02			
0.0	100.000	100.0	9.63666732959D-02	-1.66921411418D-01	1.88772520272D-02	3.33640257742D-02	34	3	5.D-17

SPHERICAL BESSEL FUNCTIONS *** ABRAMOWITZ & STEGUN PAGES 465 - 466

0.0	1.000	0.0	8.41470984808D-01	-5.40302305868D-01	-3.01168678940D-01	1.38177329068D+00			
0.0	1.000	10.0	7.11655264005D-11	-6.72215008256D+08	7.08555712150D-10	7.35887504239D+09			
0.0	1.000	30.0	5.56683126698D-43	-2.94642854750D+40	1.66916554020D-41	9.12893306467D+41			
0.0	1.000	40.0	1.53821037424D-61	-8.02845085085D+58	6.15098796870D-60	3.29064842213D+60	4	1	5.D-17
0.0	2.000	0.0	4.54648713413D-01	2.08073418274D-01	-4.3539774980D-01	3.50612004276D-01			
0.0	2.000	10.0	6.82530086497D-08	-3.55414720085D+05	3.35288171637D-07	1.91689203038D+06			
0.0	2.000	30.0	5.83661788752D-34	-1.40739387104D+31	8.73679752133D-33	2.17668398491D+32			
0.0	2.000	50.0	4.01157529034D-65	-1.23502194437D+63	1.00211458762D-64	3.1468099247D+64	4	1	5.D-17
0.0	5.000	0.0	-1.9178489492D-01	-5.6732429435D-02	9.509490792D-02	-1.80438367514D-01			
0.0	5.000	10.0	4.07344242489D-04	-2.66561144057D+01	7.21942378124D-04	5.09540067582D+01			
0.0	5.000	30.0	4.28273021730D-22	-7.76071756976D+18	2.53543820515D-21	4.74537949931D+19			
0.0	5.000	50.0	2.85747935044D-46	-6.96410918827D+42	2.84357588031D-45	7.06812704746D+43	6	1	5.D-17
0.0	10.000	0.0	-5.44021110889D-02	8.39071529076D-02	-7.84669417988D-02	-6.27928263797D-02			
0.0	10.000	10.0	6.46051544926D-02	-1.72453672088D-01	2.90307396067D-02	7.72932499058D-02			
0.0	10.000	30.0	2.51205738500D-13	-6.90831864609D+09	7.1272057762D-13	2.0207735936D+10			
0.0	10.000	50.0	2.23069602322D-31	-4.52822727235D+27	1.09348679947D-30	2.26316952204D+28	7	1	5.D-17
0.0	50.000	0.0	-5.2474977408D-03	-1.92993205698D-02	1.94042705113D-02	-4.86151066268D-03			
0.0	50.000	10.0	-1.50392214635D-02	1.35246875112D-02	-1.29162985287D-02	-1.49815666420D-02			
0.0	50.000	30.0	-1.49467345361D-03	-2.24122681205D-02	1.78070439540D-02	-6.04651449647D-04			
0.0	50.000	100.0	1.01901226293D-22	-1.12569289133D+18	1.76984656820D-22	1.97023762370D+18	13	1	5.D-17
0.0	100.000	0.0	-5.06365641100D-03	-8.62318872288D-03	8.67382528699D-03	-4.97742452387D-03			
0.0	100.000	10.0	-1.95657859714D-04	1.00257773736D-02	-9.96852414545D-03	-2.95394082628D-04			
0.0	100.000	30.0	8.70062851445D-03	-5.41292934887D-03	5.06366936035D-03	8.34315766796D-03			
0.0	100.000	100.0	1.08804770114D-02	-2.29838504916D-02	2.28730043501D-03	4.35909461714D-03	33	1	5.D-17

BESSELS WITH ORDER 1/3 *** WATSON TABLE III PAGES 714 - 729

0.0	0.240	0.3	5.46408741152D-01	-1.37617973522D+00	7.09495139391D-01	3.06764776175D+00	4	319	5.D-17
0.0	0.500	0.3	6.72830828498D-01	-8.40627826043D-01	3.19780290150D-01	1.49281941953D+00	6	160	5.D-17
0.0	1.000	0.3	7.30876402169D-01	-2.78801641276D-01	-5.52851752674D-02	8.92125353065D-01	8	83	5.D-17
0.0	2.000	0.3	4.42939818149D-01	3.43199966266D-01	-4.56138918066D-01	3.65203168163D-01	10	43	5.D-17
0.0	4.000	0.3	-3.55427373455D-01	1.79416766344D-01	-1.36039481820D-01	-3.79113116300D-01	15	22	5.D-17
0.0	5.000	0.3	-3.06420463800D-01	-1.81923211293D-01	2.12890980263D-01	-2.89126067463D-01	16	18	5.D-17
0.0	8.000	0.3	2.59776161108D-01	1.09587794634D-01	-1.25906384515D-01	2.53216724183D-01	22	12	5.D-17
0.0	10.000	0.3	-1.86145167049D-01	1.70201117883D-01	-1.61022875224D-01	-1.94771126449D-01	25	10	5.D-17
0.0	16.000	0.3	-1.04162684107D-01	1.70082756218D-01	-1.66875284009D-01	-1.09503010813D-01	33	7	5.D-17

BLANCH TEST OF RECURRENCE THROUGH THE ZERO OF J5 NEAR X = 8.7714838

0.0	8.771	0.0	-3.16762512697D-02	2.67319051367D-01	-2.65945081557D-01	-4.69170892834D-02			
0.0	8.771	1.0	2.65945081557D-01	4.69170892834D-02	-6.19955322056D-02	2.61970231218D-01			
0.0	8.771	2.0	9.23148131415D-02	-2.56621411070D-01	2.4489623894D-01	1.05429746189D-01			
0.0	8.771	4.0	-2.45434210955D-01	1.44479105276D-01	-1.11923695074D-01	-2.29828225183D-01			
0.0	8.771	5.0	3.91711840794D-09	2.95714047272D-01	-2.45434213188D-01	-2.40864726892D-02			
0.0	8.771	6.0	2.45434215421D-01	1.92652050654D-01	-1.6788535873D-01	1.63933343997D-01			
0.0	8.771	10.0	1.07825846963D-01	-6.03285674072D-01	7.11703664600D-02	2.74908871645D-01			
0.0	8.771	14.0	2.926648073665D-03	-1.00251142551D+01	3.73421850293D-03	1.20077195602D+01			
0.0	8.771	18.0	1.99004280877D-05	-1.01853380518D+03	3.59862493115D-05	1.80524423005D+03			
0.0	8.771	20.0	1.11736289101D-06	-1.58571189507D+04	2.30378879561D-06	3.22606868408D+04	10	11	5.D-17

0.0	1.000	100.0	8.45870866831D-04	-3.76330645768D+00	8.45828991019D-02	3.76308505766D+02	3	0	1.D+00
	**** IEXP =	185	F,FP *10**(-IEXP)	G,GP *10**(IEXP)					
0.0	2.000	100.0	1.06424628231D-03	-2.99155272347D+00	5.32017760004D-02	1.49546182065D+02	4	0	1.D+00
	**** IEXP =	155	F,FP *10**(-IEXP)	G,GP *10**(IEXP)					

EXAMPLES OF ERROR CONDITIONS

CF2 HAS FAILED TO CONVERGE AFTER 20000. ITERATIONS
P,Q,DP,DQ,ACCUR = 2.1651719D+02 9.2513855D+02 1.1838276D-08 -2.3027605D-08 1.000D-16

IFAIL = 2 JMAX,KFN = 1 0 0.0010 -500.0000 0.0 0.0

FOR XX = 1.000D-08 TRY SMALL-X SOLUTIONS OR X NEGATIVE
SQUARE ROOT ACCURACY PARAMETER = 1.000D-08

IFAIL = -1 JMAX,KFN = 1 0 0.0000 -500.0000 0.0 0.0

FOR XX = -2.000D+04 TRY SMALL-X SOLUTIONS OR X NEGATIVE
SQUARE ROOT ACCURACY PARAMETER = 1.000D-08

IFAIL = -1 JMAX,KFN = 1 0 -20000.0000 1.0000 0.0 -2.2000

FOR XX = -5.500D+01 TRY SMALL-X SOLUTIONS OR X NEGATIVE
SQUARE ROOT ACCURACY PARAMETER = 1.000D-08

IFAIL = -1 JMAX,KFN = 1 0 -55.0000 0.0 0.0 1.0000


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C COULFG RETURNS F,G,F',G', FOR REAL XX.GT.0,REAL ETA1 (INCLUDING 0), CABNK0069
C AND REAL LAMBDA(XLMIN) .GT. -1 FOR INTEGER-SPACED LAMBDA VALUES CABNK0070
C THUS GIVING POSITIVE-ENERGY SOLUTIONS TO THE COULOMB SCHRODINGER CABNK0071
C EQUATION, TO THE KLEIN-GORDON EQUATION AND TO SUITABLE FORMS OF CABNK0072
C THE DIRAC EQUATION ,ALSO SPHERICAL & CYLINDRICAL BESSEL EQUATIONS CABNK0073
C CABNK0074
C FOR A RANGE OF LAMBDA VALUES (XLMAX - XLMIN) MUST BE AN INTEGER, CABNK0075
C STARTING ARRAY ELEMENT IS M1 = MAXO(IDINT(XLMIN+ACCUR),0) + 1 CABNK0076
C SEE TEXT FOR MODIFICATIONS FOR INTEGER L-VALUES CABNK0077
C CABNK0078
C IF 'MODE' = 1 GET F,G,F',G' FOR INTEGER-SPACED LAMBDA VALUES CABNK0079
C = 2 F,G UNUSED ARRAYS MUST BE DIMENSIONED IN CABNK0080
C = 3 F CALL TO AT LEAST LENGTH (1) CABNK0081
C IF 'KFN' = 0 REAL COULOMB FUNCTIONS ARE RETURNED CABNK0082
C = 1 SPHERICAL BESSEL " " " CABNK0083
C = 2 CYLINDRICAL BESSEL " " " CABNK0084
C THE USE OF 'MODE' AND 'KFN' IS INDEPENDENT CABNK0085
C CABNK0086
C PRECISION: RESULTS TO WITHIN 2-3 DECIMALS OF 'MACHINE ACCURACY' CABNK0087
C IN OSCILLATING REGION X .GE. ETA1 + SQRT(ETA1**2 + XLM(XLM+1)) CABNK0088
C COULFG IS CODED FOR REAL*8 ON IBM OR EQUIVALENT ACCUR = 10**-16 CABNK0089
C USE AUTODBL + EXTENDED PRECISION ON HX COMPILER ACCUR = 10**-33 CABNK0090
C FOR MANPISAS OF 56 & 112 BITS. FOR SINGLE PRECISION CDC (48 BITS) CABNK0091
C REASSIGN DSQRT=SQRT ETC. SEE TEXT FOR COMPLEX ARITHMETIC VERSION CABNK0092
C ***** CABNK0093
C IMPLICIT REAL*8 (A-H,O-Z) ABNK0094
C DIMENSION FC(1),GC(1),FCP(1),GCP(1) ABNK0095
C LOGICAL ETANEO,XLTURN ABNK0096
C COMMON /STEED/ PACCQ,NFP,NPQ,IEXP,M1 ABNK0097
C *** COMMON BLOCK IS FOR INFORMATION ONLY. NOT REQUIRED IN CODE ABNK0099
C *** COULFG HAS CALLS TO: DSQRT,DABS,DMOD,IDINT,DSIGN,DFLOAT,DMIN1 ABNK0100
C DATA ZERO,ONE,TWO,TEN2,ABORT /0.0D0, 1.0D0, 2.0D0, 1.0D2, 2.0D4/ ABNK0101
C DATA HALF,TM30 / 0.5D0, 1.0D-30 / ABNK0102
C DATA RT2DPI /0.79788 45608 02865 35587 98921 19868 76373 D0/ ABNK0103
C *** THIS CONSTANT IS DSQRT(TWO/PI): USE QO FOR IBM REAL*16: DO FOR ABNK0104
C *** REAL*8 & CDC DOUBLE P: EO FOR CDC SINGLE P; AND TRUNCATE VALUE. ABNK0105
C ABNK0106
C ACCUR = 1.0D-16 ABNK0107
C *** CHANGE ACCUR TO SUIT MACHINE AND PRECISION REQUIRED ABNK0108
C MODE = 1 ABNK0109
C IF(MODE1 .EQ. 2 .OR. MODE1 .EQ. 3 ) MODE = MODE1 ABNK0110
C IFAIL = 0 ABNK0111
C IEXP = 1 ABNK0112
C NPQ = 0 ABNK0113
C ETA = ETA1 ABNK0114
C GJWKB = ZERO ABNK0115
C PACCQ = ONE ABNK0116
C IF(KFN .NE. 0) ETA = ZERO ABNK0117
C ETANEO = ETA .NE. ZERO ABNK0118
C ACC = ACCUR ABNK0119
C ACC4 = ACC*TEN2*TEN2 ABNK0120
C ACCH = DSQRT(ACC) ABNK0121
C *** TEST RANGE OF XX, EXIT IF .LE. DSQRT(ACCUR) OR IF NEGATIVE ABNK0122
C ABNK0123
C IF(XX .LE. ACCH) GO TO 100 ABNK0124
C X = XX ABNK0125
C XLM = XLMIN ABNK0126
C IF(KFN .EQ. 2) XLM = XLM - HALF ABNK0127
C IF(XLM .LE. -ONE .OR. XLMAX .LT. XLMIN) GO TO 105 ABNK0128
C E2MML = ETA*ETA + XLM*XLM + XLM ABNK0129
C XLTURN = X*(X - TWO*ETA) .LT. XLM*XLM + XLM ABNK0130
C DELL = XLMAX - XLMIN + ACC ABNK0131
C IF(DABS(DMOD(DELL,ONE)) .GT. ACC) WRITE(6,2040)XLMAX,XLMIN,DELL ABNK0132
C LXTRA = IDINT(DELL) ABNK0133
C XLL = XLM + DFLOAT(LXTRA) ABNK0134
C *** LXTRA IS NUMBER OF ADDITIONAL LAMBDA VALUES TO BE COMPUTED ABNK0135
C *** XLL IS MAX LAMBDA VALUE, OR 0.5 SMALLER FOR J,Y BESSELS ABNK0136
C *** DETERMINE STARTING ARRAY ELEMENT (M1) FROM XLMIN ABNK0137
C M1 = MAXO(IDINT(XLMIN + ACC),0) + 1 ABNK0138
C L1 = M1 + LXTRA ABNK0139
C ABNK0140

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C *** EVALUATE CF1 = F = FPRIME(XL,ETA,X)/F(XL,ETA,X) ABNK0141
C ABNK0142
  XI = ONE/X ABNK0143
  FCL = ONE ABNK0144
  PK = XLL + ONE ABNK0145
  PX = PK + ABORT ABNK0146
  2 EK = ETA / PK ABNK0147
  F = (EK + PK*XI)*FCL + (FCL - ONE)*XI ABNK0148
  PK1 = PK + ONE ABNK0149
C *** TEST ENSURES BL.NE. ZERO FOR NEGATIVE ETA; FIXUP IS EXACT. ABNK0150
  IF(DABS(ETA*X + PK*PK1) .GT. ACC) GO TO 3 ABNK0151
  FCL = (ONE + EK*EK)/(ONE + (ETA/PK1)**2) ABNK0152
  PK = TWO + PK ABNK0153
  GO TO 2 ABNK0154
  3 D = ONE/((PK + PK1)*(XI + EK/PK1)) ABNK0155
  DF = -FCL*(ONE + EK*EK)*D ABNK0156
  IF(FCL .NE. ONE) FCL = -ONE ABNK0157
  IF(D .LT. ZERO) FCL = -FCL ABNK0158
  F = F + DF ABNK0159
C ABNK0160
C *** BEGIN CF1 LOOP ON PK = K = LAMBDA + 1 ABNK0161
C ABNK0162
  P = ONE ABNK0163
  4 PK = PK1 ABNK0164
  PK1 = PK1 + ONE ABNK0165
  EK = ETA / PK ABNK0166
  TK = (PK + PK1)*(XI + EK/PK1) ABNK0167
  D = TK - D*(ONE + EK*EK) ABNK0168
  IF(DABS(D) .GT. ACCH) GO TO 5 ABNK0169
  WRITE (6,1000) D,DF,ACCH,PK,EK,ETA,X ABNK0170
  P = P + ONE ABNK0171
  IF( P .GT. TWO ) GO TO 110 ABNK0172
  5 D = ONE/D ABNK0173
  IF( D .LT. ZERO) FCL = -FCL ABNK0174
  DF = DF*(D*TK - ONE) ABNK0175
  F = F + DF ABNK0176
  IF(PK .GT. PX) GO TO 110 ABNK0177
  IF(DABS(DF) .GE. DABS(F)*ACC) GO TO 4 ABNK0178
  NFP = PK - XLL - 1 ABNK0179
  IF(LXTRA .EQ. 0) GO TO 7 ABNK0180
C ABNK0181
C *** DOWNWARD RECURRENCE TO LAMBDA = XLM. ARRAY GC, IF PRESENT, STORES RL ABNK0182
C ABNK0183
  FCL = FCL*TM30 ABNK0184
  FPL = FCL*F ABNK0185
  IF(MODE .EQ. 1) FCP(L1) = FPL ABNK0186
  FC(L1) = FCL ABNK0187
  XL = XLL ABNK0188
  RL = ONE ABNK0189
  EL = ZERO ABNK0190
  DO 6 LP = 1, LXTRA ABNK0191
  IF(ETANEO) EL = ETA/XL ABNK0192
  IF(ETANEO) RL = DSQRT(ONE + EL*EL) ABNK0193
  SL = EL + XL*XI ABNK0194
  L = L1 - LP ABNK0195
  FCL1 = (FCL *SL + FPL)/RL ABNK0196
  FPL = FCL1*SL - FCL *RL ABNK0197
  FCL = FCL1 ABNK0198
  FC(L) = FCL ABNK0199
  IF(MODE .EQ. 1) FCP(L) = FPL ABNK0200
  IF(MODE .NE. 3 .AND. ETANEO) GC(L+1) = RL ABNK0201
  6 XL = XL - ONE ABNK0202
  IF(FCL .EQ. ZERO) FCL = ACC ABNK0203
  F = FPL/FCL ABNK0204
C *** NOW WE HAVE REACHED LAMBDA = XLMIN = XLM ABNK0205
C *** EVALUATE CF2 = P + I.Q AGAIN USING STEED'S ALGORITHM ABNK0206
C *** SEE TEXT FOR COMPACT COMPLEX CODE FOR SP CDC OR NON-ANSI IBM ABNK0207
C ABNK0208
  7 IF( XLTURN ) CALL JWKB(X,ETA,DMAX1(XLM,ZERO),FJWKB,GJWKB,IEXP) ABNK0209
  IF( IEXP .GT. 1 .OR. GJWKB .GT. ONE/(ACCH*TEN2)) GO TO 9 ABNK0210
  XLTURN = .FALSE. ABNK0211
  TA = TWO*ABORT ABNK0212
  PK = ZERO ABNK0213
  WI = ETA + ETA ABNK0214

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P = ZERO
Q = ONE - ETA*XI
AR = -E2MM1
AI = ETA
BR = TWO*(X - ETA)
BI = TWO
DR = BR/(BR*BR + BI*BI)
DI = -BI/(BR*BR + BI*BI)
DP = -XI*(AR*DI + AI*DR)
DQ = XI*(AR*DR - AI*DI)
8 P = P + DP
Q = Q + DQ
PK = PK + TWO
AR = AR + PK
AI = AI + WI
BI = BI + TWO
D = AR*DR - AI*DI + BR
DI = AI*DR + AR*DI + BI
C = ONE/(D*D + DI*DI)
DR = C*D
DI = -C*DI
A = BR*DR - BI*DI - ONE
B = BI*DR + BR*DI
C = DP*A - DQ*B
DQ = DP*B + DQ*A
DP = C
IF(PK .GT. TA) GO TO 120
IF(DABS(DP)+DABS(DQ).GE.(DABS(P)+DABS(Q))*ACC) GO TO 8
NPQ = PK/TWO
PACCQ = HALF*ACC/DMIN1(DABS(Q),ONE)
IF(DABS(P) .GT. DABS(Q)) PACCQ = PACCQ*DABS(P)
C
C *** SOLVE FOR FCM = F AT LAMBDA = XLM, THEN FIND NORM FACTOR W=W/FCM
C
GAM = (F - P)/Q
IF(Q .LE. ACC4*DABS(P)) GO TO 130
W = ONE/DSQRT((F - P)*GAM + Q)
GO TO 10
C *** ARRIVE HERE IF G(XLM) .GT. 10**6 OR IEXP .GT. 70 & XLTURN = .TRUE.
9 W = FJWKB
GAM = GJWKB*W
P = F
Q = ONE
C
C *** NORMALISE FOR SPHERICAL OR CYLINDRICAL BESSEL FUNCTIONS
C
10 ALPHA = ZERO
IF(KFN .EQ. 1) ALPHA = XI
IF(KFN .EQ. 2) ALPHA = XI*HALF
BETA = ONE
IF(KFN .EQ. 1) BETA = XI
IF(KFN .EQ. 2) BETA = DSQRT(XI)*RT2DPI
FCM = DSIGN(W,FCL)*BETA
FC(M1) = FCM
IF(MODE .EQ. 3) GO TO 11
IF(.NOT. XLTURN) GCL = FCM*GAM
IF(XLTURN) GCL = GJWKB*BETA
IF(KFN .NE. 0) GCL = -GCL
GC(M1) = GCL
GPL = GCL*(P - Q/GAM) - ALPHA*GCL
IF(MODE .EQ. 2) GO TO 11
GCP(M1) = GPL
FCP(M1) = FCM*(F - ALPHA)
11 IF(LXTRA .EQ. 0) RETURN
C *** UPWARD RECURRENCE FROM GC(M1),GCP(M1) STORED VALUE IS RL
C *** RENORMALISE FC,FCP AT EACH LAMBDA AND CORRECT REGULAR DERIVATIVE
C *** XL = XLM HERE AND RL = ONE, EL = ZERO FOR BESSELS
W = BETA*W/DABS(FCL)
MAXL = L1 + 1
DO 12 L = M1,MAXL
IF(MODE .EQ. 3) GO TO 12

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                XL = XL + ONE                ABNK0286
IF(ETANEO)      EL = ETA/XL                ABNK0287
IF(ETANEO)      RL = GC(L+1)              ABNK0288
                SL = EL + XL*XI            ABNK0289
GCL1            = ((SL - ALPHA)*GCL - GPL)/RL ABNK0290
GPL             = RL*GCL - (SL + ALPHA)*GCL1 ABNK0291
GCL            = GCL1                      ABNK0292
GC(L+1)        = GCL1                      ABNK0293
                IF(MODE .EQ. 2)           GO TO 12 ABNK0294
GCP(L+1)        = GPL                      ABNK0295
FCP(L+1)        = W*(FCP(L+1) - ALPHA*FC(L+1)) ABNK0296
12 FC(L+1)      = W* FC(L+1)              ABNK0297
RETURN                                                 ABNK0298
1000 FORMAT(/' CF1 ACCURACY LOSS: D,DF,ACCH,K,ETA/K,ETA,X = ',1P7D9.2/) ABNK0299
C
C ***      ERROR MESSAGES
C
100 IFAIL = -1
WRITE(6,2000) XX,ACCH
2000 FORMAT(' FOR XX = ',1PD12.3,' TRY SMALL-X SOLUTIONS',
*' OR X NEGATIVE'/' , ' SQUARE ROOT ACCURACY PARAMETER = ',D12.3/)
RETURN
105 IFAIL = -2
WRITE (6,2005) XLMAX,XLMIN,XLM
2005 FORMAT(/' PROBLEM WITH INPUT ORDER VALUES:XLMAX,XLMIN,XLM = ',
*'1P3D15.6/)
RETURN
110 IFAIL = 1
WRITE (6,2010) ABORT,F ,DF,PK,PX,ACC
2010 FORMAT(' CF1 HAS FAILED TO CONVERGE AFTER ',F10.0,' ITERATIONS',/
*' F,DF,PK,PX,ACCUR = ',1P5D12.3//)
RETURN
120 IFAIL = 2
WRITE (6,2020) ABORT,P,Q,DP,DQ,ACC
2020 FORMAT(' CF2 HAS FAILED TO CONVERGE AFTER ',F7.0,' ITERATIONS',/
*' P,Q,DP,DQ,ACCUR = ',1P4D17.7,D12.3//)
RETURN
130 IFAIL = 3
WRITE (6,2030) P,Q,ACC,DELL,LXTRA,M1
2030 FORMAT(' FINAL Q.LE.DABS(P)*ACC*10**4 , P,Q,ACC = ',1P3D12.3,4X,
*' DELL,LXTRA,M1 = ',D12.3,2I5 /)
RETURN
2040 FORMAT(' XLMAX - XLMIN = DELL NOT AN INTEGER ',1P3D20.10/)
END
C
SUBROUTINE JWKB(XX,ETA1,XL,FJWKB,GJWKB,IEXP)
REAL*8      XX,ETA1,XL,FJWKB,GJWKB,DZERO
C *** COMPUTES JWKB APPROXIMATIONS TO COULOMB FUNCTIONS FOR XL.GE. 0 ABNK0333
C *** AS MODIFIED BY BIEDENHARN ET AL. PHYS REV 97 (1955) 542-554 ABNK0334
C *** CALLS DMAX1,SQRT,ALOG,EXP,ATAN2,FLOAT,INT BARNETT FEB 1981 ABNK0335
DATA ZERO,HALF,ONE,SIX,TEN/ 0.0EO, 0.5EO, 1.0EO, 6.0EO, 10.0EO /ABNK0336
DATA DZERO, RL35, ALOGE /0.0DO, 35.0EO, 0.43429 45 EO / ABNK0337
X = XX
ETA = ETA1
GH2 = X*(ETA + ETA - X)
XLL1 = DMAX1(XL*XL + XL,DZERO)
IF(GH2 + XLL1 .LE. ZERO) RETURN
HLL = XLL1 + SIX/RL35
HL = SQRT(HLL)
SL = ETA/HL + HL/X
RL2 = ONE + ETA*ETA/HLL
GH = SQRT(GH2 + HLL)/X
PHI = X*GH - HALF*( HL*ALOG((GH + SL)**2/RL2) - ALOG(GH) )
IF(ETA .NE. ZERO) PHI = PHI - ETA*ATAN2(X*GH,X - ETA)
PHI10 = -PHI*ALOG
IEXP = INT(PHI10)
IF(IEXP .GT. 70) GJWKB = TEN**(PHI10 - FLOAT(IEXP))
IF(IEXP .LE. 70) GJWKB = EXP(-PHI)
IF(IEXP .LE. 70) IEXP = 0
FJWKB = HALF/(GH*GJWKB)
RETURN
END
//G.SYSIN DD *

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